

Null Geodesics in de-Sitter Universe

J.A.N.K. Jayakody¹, L.N.K. de Silva²

Department of Physics¹, University of Kelaniya

Department of Mathematics² University of Kelaniya

ABSTRACT

Consider the de-Sitter universe given by

$$ds^2 = c^2 dt^2 - e^{2Ht} \left[dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right]^{[1]}, \quad (1)$$

which involves the cosmological constant.

Here H is the Hubble constant and $H = \sqrt{\frac{\Lambda}{3}}$; Λ is the cosmological constant.

For null-geodesics $ds^2 = 0$.

$$\text{Let } d\mu^2 = c^2 dt^2 - e^{2Ht} \left[dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right]. \quad (2)$$

Dividing both sides of (2) by the element $d\mu^2$,

$$L = c^2 \left(\frac{dt}{d\mu} \right)^2 - e^{2Ht} \left(\frac{dr}{d\mu} \right)^2 - e^{2Ht} r^2 \left[\left(\frac{d\theta}{d\mu} \right)^2 + \sin^2 \theta \left(\frac{d\phi}{d\mu} \right)^2 \right] = 0. \quad (3)$$

The equations of the null-geodesics can be written by using the Euler-Lagrange's equations

$$\frac{d}{d\mu} \left(\frac{\partial L}{\partial \dot{x}} \right) - \frac{\partial L}{\partial x} = 0, \quad (4)$$

where $x = (t, r, \theta, \phi)$ and dot denotes differentiation with respect to μ .

Considering the t coordinate,

$$\frac{d}{d\mu} \left[2c^2 \left(\frac{dt}{d\mu} \right) \right] + 2He^{2Ht} \left(\frac{dr}{d\mu} \right)^2 + 2He^{2Ht} r^2 \left[\left(\frac{d\theta}{d\mu} \right)^2 + \sin^2 \theta \left(\frac{d\phi}{d\mu} \right)^2 \right] = 0. \quad (5)$$

Considering the θ coordinate,

$$\frac{d}{d\mu} \left[-2e^{2Ht} r^2 \left(\frac{d\theta}{d\mu} \right) \right] + 2e^{2Ht} r^2 \sin \theta \cos \theta \left(\frac{d\phi}{d\mu} \right)^2 = 0. \quad (6)$$

Considering the ϕ coordinate,

$$\frac{d}{d\mu} \left[-2e^{2Ht} r^2 \sin^2 \theta \left(\frac{d\phi}{d\mu} \right) \right] = 0. \quad (7)$$

Since $ds = 0$, for null-geodesics we also have

$$c^2 \left(\frac{dt}{d\mu} \right)^2 - e^{2Ht} \left(\frac{dr}{d\mu} \right)^2 - e^{2Ht} r^2 \left[\left(\frac{d\theta}{d\mu} \right)^2 + \sin^2 \theta \left(\frac{d\phi}{d\mu} \right)^2 \right] = 0 \quad (8)$$

These are the null-geodesics equations of the de-Sitter space-time.

Equation (6) is satisfied by $\theta = \frac{\pi}{2}$.

From (5),

$$\frac{d}{d\mu} \left[2c^2 \left(\frac{dt}{d\mu} \right) \right] + 2He^{2Ht} \left(\frac{dr}{d\mu} \right)^2 + 2He^{2Ht} r^2 \left(\frac{d\phi}{d\mu} \right)^2 = 0 \quad (9)$$

From (7),

$$\begin{aligned} \frac{d}{d\mu} \left[-2e^{2Ht} r^2 \left(\frac{d\phi}{d\mu} \right) \right] &= 0 \\ -2e^{2Ht} r^2 \left(\frac{d\phi}{d\mu} \right) &= c_1 \quad ; \quad \text{where } c_1 \text{ is a constant} \end{aligned}$$

$$d\mu = \frac{-2e^{2Ht} r^2}{c_1} d\phi \quad (10)$$

From (8),

$$c^2 \left(\frac{dt}{d\mu} \right)^2 - e^{2Ht} \left(\frac{dr}{d\mu} \right)^2 - e^{2Ht} r^2 \left(\frac{d\phi}{d\mu} \right)^2 = 0 \quad (11)$$

From (9) and (10), we obtain,

$$\left(\frac{dr}{d\phi} \right)^2 = - \left(\frac{d^2 t}{d\mu^2} + \frac{Hc_1^2}{4c^2 e^{2Ht} r^2} \right) \frac{4c^2 e^{2Ht} r^4}{Hc_1^2} \quad (12)$$

From (10) and (11), we have,

$$\left(\frac{dr}{d\phi}\right)^2 = \left[\left(\frac{dt}{d\mu}\right)^2 - \frac{c_1^2}{4c^2 e^{2Ht} r^2}\right] \frac{4c^2 e^{2Ht} r^4}{c_1^2} \quad (13)$$

(12) and (13) give,

$$\frac{1}{H} \left(\frac{d^2 t}{d\mu^2}\right) + \left(\frac{dt}{d\mu}\right)^2 = 0$$

Let $\frac{dt}{d\mu} = p$. Then $\frac{d^2 t}{d\mu^2} = p \frac{dp}{dt}$

Let $p^2 = v$. Then $2p \frac{dp}{dt} = \frac{dv}{dt}$ and finally we have,

$$\frac{dt}{d\mu} = \sqrt{e^{-(2Ht+c_2)}} \quad ; \quad \text{where } c_2 \text{ is a constant} \quad (14)$$

From (13) and (14), substituting $r = \frac{1}{u}$ we obtain

$$\frac{du}{d\phi} = \sqrt{\frac{4c^2 e^{-c_2}}{c_1^2} - u^2}$$

$$\sin^{-1} \frac{u}{\sqrt{\frac{4c^2 e^{-c_2}}{c_1^2}}} = \phi + c_3 \quad ; \quad \text{where } c_3 \text{ is a constant.}$$

$$u = \frac{2c}{c_1} \sqrt{e^{-c_2}} \sin(\phi + c_3)$$

Thus the equation of a null geodesic in de Sitter universe is given by the above expression

which agrees with the equation $u = \left(\frac{E^2}{h^2} + \frac{\Lambda}{3}\right)^{\frac{1}{2}} \sin(\phi + A)$ [2] we had in the case of Schwarzschild - de Sitter metric with $m = 0$, after identifying the corresponding constants.

References

[1] Narlikar, J.V., (1993) *Introduction to Cosmology*, Cambridge University Press.

[2] Jayakody, J.A.N.K., de Silva, L.N.K. *Path of a light ray near a body with cosmological constant*, 10th Annual Research Symposium 2009, University of Kelaniya (2009).